### Non-scattering Wave Numbers

Michael Vogelius

Department of Mathematics Rutgers University

Joint with Jingni Xiao

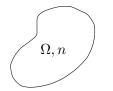
(ロ)、(型)、(E)、(E)、 E) の(()

#### Introduction

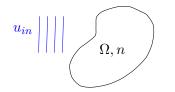
Finiteness of Non-scattering Wave Numbers

Connection with Schiffer's conjecture

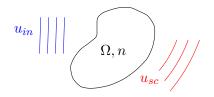
◆□▶ ◆□▶ ◆ 臣▶ ◆ 臣▶ ○ 臣 ○ の Q @



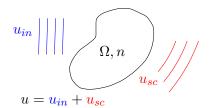
◆□▶ ◆□▶ ◆ 臣▶ ◆ 臣▶ ○ 臣 ○ の Q @





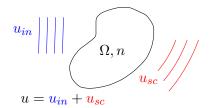






イロト イヨト イヨト イヨト

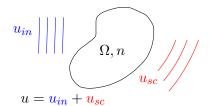
æ



$$\begin{split} \Delta u + k^2 n u &= 0 \text{ in } \mathbb{R}^d \\ n &= 1 \text{ in } \Omega^c \\ n &= n_0 \neq 1 \text{ in } \Omega \end{split}$$

ヘロト 人間ト 人間ト 人間ト

æ



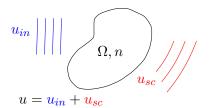
$$\begin{split} \Delta u + k^2 n u &= 0 \text{ in } \mathbb{R}^d \\ n &= 1 \text{ in } \Omega^c \\ n &= n_0 \neq 1 \text{ in } \Omega \end{split}$$

・ロト ・ 国 ト ・ ヨ ト ・ ヨ ト

3

 $u_{in} = u_{k,in}$ : incident wave

$$\Delta u_{k,in} + k^2 u_{k,in} = 0 \quad \text{in } \mathbb{R}^d;$$



$$\begin{split} \Delta u + k^2 n u &= 0 \text{ in } \mathbb{R}^d \\ n &= 1 \text{ in } \Omega^c \\ n &= n_0 \neq 1 \text{ in } \Omega \end{split}$$

 $u_{in} = u_{k,in}$ : incident wave

$$\Delta u_{k,in} + k^2 u_{k,in} = 0$$
 in  $\mathbb{R}^d$ ;

 $u_{sc} = u_{k,sc}$ : scattered wave, with Sommerfeld radiation condition

$$\lim_{r \to \infty} r^{\frac{d-1}{2}} \left( \frac{\partial}{\partial r} u_{k,sc} - iku_{k,sc} \right) = 0 \quad \text{uniformly for all directions.}$$

Suppose there exits  $(k, u_{k,in})$  such that  $u_{k,sc} = 0$  in  $\Omega^c$ .

◆□▶ ◆□▶ ◆ 臣▶ ◆ 臣▶ ○ 臣 ○ の Q @

Suppose there exits  $(k, u_{k,in})$  such that  $u_{k,sc} = 0$  in  $\Omega^c$ . Then

$$\begin{split} \Delta u_{k,in} + k^2 u_{k,in} &= 0 & \text{in } \Omega, \\ \Delta u_{k,sc} + k^2 n_0 u_{k,sc} &= k^2 (1 - n_0) u_{k,in} & \\ u_{k,sc} &= \partial_\nu u_{k,sc} &= 0 & \text{on } \partial \Omega. \end{split}$$

(ロ)、(型)、(E)、(E)、 E) の(()

Suppose there exits  $(k, u_{k,in})$  such that  $u_{k,sc} = 0$  in  $\Omega^c$ . Then

$$\begin{split} \Delta u_{k,in} + k^2 u_{k,in} &= 0 & \text{in } \Omega, \\ \Delta u_{k,sc} + k^2 n_0 u_{k,sc} &= k^2 (1 - n_0) u_{k,in} & \\ u_{k,sc} &= \partial_{\nu} u_{k,sc} &= 0 & \text{on } \partial \Omega. \end{split}$$

— Alternatively  $u_{k,sc}$  satisfies

$$\begin{split} (\Delta+k^2)\frac{1}{n_0-1}(\Delta+k^2n_0)u_{k,sc} &= 0 \mbox{ in } \Omega \\ \mbox{with} \qquad u_{k,sc} &= \partial_\nu u_{k,sc} = 0 \mbox{ on } \partial\Omega \end{split}$$

▲□▶ ▲□▶ ▲ 三▶ ▲ 三▶ 三三 - のへぐ

Suppose there exits  $(k, u_{k,in})$  such that  $u_{k,sc} = 0$  in  $\Omega^c$ . Then

$$\begin{split} \Delta u_{k,in} + k^2 u_{k,in} &= 0 & \text{in } \Omega_i \\ \Delta u_{k,sc} + k^2 n_0 u_{k,sc} &= k^2 (1 - n_0) u_{k,in} & \\ u_{k,sc} &= \partial_\nu u_{k,sc} &= 0 & \text{on } \partial \Omega. \end{split}$$

— Alternatively  $u_{k,sc}$  satisfies

$$\begin{split} (\Delta+k^2)\frac{1}{n_0-1}(\Delta+k^2n_0)u_{k,sc} &= 0 \mbox{ in } \Omega \\ \mbox{with} \qquad u_{k,sc} &= \partial_\nu u_{k,sc} = 0 \mbox{ on } \partial\Omega \end{split}$$

▲□▶ ▲□▶ ▲ 三▶ ▲ 三▶ 三 のへぐ

- the interior transmission eigenvalue problem (ITEP)

Suppose there exits  $(k, u_{k,in})$  such that  $u_{k,sc} = 0$  in  $\Omega^c$ . Then

$$\begin{aligned} \Delta u_{k,in} + k^2 u_{k,in} &= 0 & \text{in } \Omega, \\ \Delta u_{k,sc} + k^2 n_0 u_{k,sc} &= k^2 (1 - n_0) u_{k,in} & \\ u_{k,sc} &= \partial_{\nu} u_{k,sc} &= 0 & \text{on } \partial\Omega. \end{aligned}$$

— Alternatively  $u_{k,sc}$  satisfies

$$\begin{split} (\Delta+k^2)\frac{1}{n_0-1}(\Delta+k^2n_0)u_{k,sc} &= 0 \text{ in } \Omega\\ \text{with} \qquad u_{k,sc} &= \partial_\nu u_{k,sc} = 0 \quad \text{ on } \partial\Omega \end{split}$$

— the interior transmission eigenvalue problem (ITEP) and  $u_{k,in}$  is then defined by

$$u_{k,in} = \frac{1}{k^2(1-n_0)} (\Delta + k^2 n_0) u_{k,sc} .$$

Suppose there exits  $(k, u_{k,in})$  such that  $u_{k,sc} = 0$  in  $\Omega^c$ . Then

$$\begin{split} \Delta u_{k,in} + k^2 u_{k,in} &= 0 & \text{in } \Omega_i \\ \Delta u_{k,sc} + k^2 n_0 u_{k,sc} &= k^2 (1 - n_0) u_{k,in} & \\ u_{k,sc} &= \partial_\nu u_{k,sc} &= 0 & \text{on } \partial \Omega. \end{split}$$

— Alternatively  $u_{k,sc}$  satisfies

$$\begin{split} (\Delta+k^2)\frac{1}{n_0-1}(\Delta+k^2n_0)u_{k,sc} &= 0 \text{ in } \Omega\\ \text{with} \qquad u_{k,sc} &= \partial_\nu u_{k,sc} = 0 \quad \text{ on } \partial\Omega \end{split}$$

— the interior transmission eigenvalue problem (ITEP) and  $u_{k,in}$  is then defined by

$$u_{k,in} = \frac{1}{k^2(1-n_0)} (\Delta + k^2 n_0) u_{k,sc} .$$

- Notice: this (ITEP) is only a **necessary** condition for Non-scattering. We restrict our attention to real wave numbers k.

► Observation:

 $\mathscr{A} := \{\mathsf{non-scattering wave numbers}\} \subset \{\mathsf{ITEV}\} := \mathscr{B}$ 

We restrict our attention to real wave numbers k.

Observation:

 $\mathscr{A} := \{\mathsf{non-scattering wave numbers}\} \subset \{\mathsf{ITEV}\} := \mathscr{B}$ 



$$#\mathscr{B} = \infty$$
(countable) cf., [Cakoni-Colton-Haddar (book)]

We restrict our attention to real wave numbers k.

Observation:

 $\mathscr{A} := \{ \mathsf{non-scattering wave numbers} \} \subset \{ \mathsf{ITEV} \} := \mathscr{B}$ 



 $#\mathscr{B} = \infty$ (countable) cf., [Cakoni-Colton-Haddar (book)]



$$\mathscr{A} = \mathscr{B} ? \qquad \# \mathscr{A} = \infty ?$$

・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・

\$\mathcal{A} = \mathcal{B}\$: (countably) infinitely many non-scattering wave numbers.

- \$\mathcal{A} = \mathcal{B}\$: (countably) infinitely many non-scattering wave numbers.
  - Spherically stratified media n = n(r) [Colton-Monk '88] Eigenfunctions/non-scattering incident waves: Herglotz functions

▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□ ● ● ●

- \$\mathcal{A} = \mathcal{B}\$: (countably) infinitely many non-scattering wave numbers.
  - Spherically stratified media n = n(r) [Colton-Monk '88] Eigenfunctions/non-scattering incident waves: Herglotz functions

A = Ø ⊊ B: no non-scattering wave numbers
 Media (domains) with singularities at the shape:

- \$\mathcal{A} = \mathcal{B}\$: (countably) infinitely many non-scattering wave numbers.
  - Spherically stratified media n = n(r) [Colton-Monk '88] Eigenfunctions/non-scattering incident waves: Herglotz functions
- A = Ø ⊊ B: no non-scattering wave numbers
  Media (domains) with singularities at the shape:

 corner, edge
 [Blåsten, Päivärinta, Sylvester, Salo, Vesalainen, Elschner, Hu, Liu, Xiao, Cakoni, Cao, Diao, etc.]

- \$\mathcal{A} = \mathcal{B}\$: (countably) infinitely many non-scattering wave numbers.
  - Spherically stratified media n = n(r) [Colton-Monk '88] Eigenfunctions/non-scattering incident waves: Herglotz functions
- A = Ø ⊊ B: no non-scattering wave numbers
  Media (domains) with singularities at the shape:

 corner, edge
 [Blåsten, Päivärinta, Sylvester, Salo, Vesalainen, Elschner, Hu, Liu, Xiao, Cakoni, Cao, Diao, etc.]

- A = B: (countably) infinitely many non-scattering wave
  numbers.
  - Spherically stratified media n = n(r) [Colton-Monk '88] Eigenfunctions/non-scattering incident waves: Herglotz functions
- A = Ø ⊊ B: no non-scattering wave numbers
  Media (domains) with singularities at the shape:
  - corner, edge
    [Blåsten, Päivärinta, Sylvester, Salo, Vesalainen, Elschner, Hu, Liu, Xiao, Cakoni, Cao, Diao, etc.]

— Let us, as an example, consider the case  $\Omega = B_1$  (ball of radius 1) in  $\mathbb{R}^2$  with  $n_0$  a constant.

We take 
$$u_{in}(x) = e^{-im\theta} J_m(kr)$$
.

◆□▶ ◆□▶ ◆三▶ ◆三▶ ◆□▶

We take  $u_{in}(x)=e^{-im\theta}J_m(kr).$  This solves  $(\Delta+k^2)u_{in}=0,$  provided  $J_m$  is the (bounded) solution to

$$\left[r^2\frac{d^2}{dr^2} + r\frac{d}{dr} + (r^2 - m^2)\right]J_m(r) = 0.$$
 BESSEL'S EQUATION

▲□▶ ▲□▶ ▲ 三▶ ▲ 三▶ 三三 - のへぐ

We take  $u_{in}(x) = e^{-im\theta}J_m(kr)$ . This solves  $(\Delta + k^2)u_{in} = 0$ , provided  $J_m$  is the (bounded) solution to

$$\left[r^2\frac{d^2}{dr^2} + r\frac{d}{dr} + (r^2 - m^2)\right]J_m(r) = 0.$$
 BESSEL'S EQUATION

It corresponds to  $u_{sc}=0$  in  $\Omega^c$  (and thus also to an ITEV) iff we can find a solution to

$$(\Delta + k^2 n(x))u = 0$$
 in  $\mathbb{R}^2$ , with  $u = e^{-im\theta} J_m(kr)$  in  $\Omega^c$ .

▲□▶▲□▶▲≡▶▲≡▶ ≡ めぬぐ

We take  $u_{in}(x) = e^{-im\theta}J_m(kr)$ . This solves  $(\Delta + k^2)u_{in} = 0$ , provided  $J_m$  is the (bounded) solution to

$$\left[r^2 \frac{d^2}{dr^2} + r \frac{d}{dr} + (r^2 - m^2)\right] J_m(r) = 0 . \quad \text{BESSEL'S EQUATION}$$

It corresponds to  $u_{sc}=0$  in  $\Omega^c$  (and thus also to an ITEV) iff we can find a solution to

$$(\Delta + k^2 n(x))u = 0$$
 in  $\mathbb{R}^2$ , with  $u = e^{-im\theta} J_m(kr)$  in  $\Omega^c$ .

This solution must have the form  $u(x)=Ae^{-im\theta}J_m(k\sqrt{n_0}r)$  inside  $\Omega,$  and so

We take  $u_{in}(x) = e^{-im\theta}J_m(kr)$ . This solves  $(\Delta + k^2)u_{in} = 0$ , provided  $J_m$  is the (bounded) solution to

$$\left[r^2 \frac{d^2}{dr^2} + r \frac{d}{dr} + (r^2 - m^2)\right] J_m(r) = 0.$$
 BESSEL'S EQUATION

It corresponds to  $u_{sc}=0$  in  $\Omega^c$  (and thus also to an ITEV) iff we can find a solution to

$$(\Delta+k^2n(x))u=0 \ \, \text{in} \ \, \mathbb{R}^2 \ , \quad \text{with} \ \, u=e^{-im\theta}J_m(kr) \ \, \text{in} \ \, \Omega^c \ .$$

This solution must have the form  $u(x) = Ae^{-im\theta}J_m(k\sqrt{n_0}r)$ inside  $\Omega$ , and so we have  $u_{sc} = 0$  in  $\Omega^c$  (and thus also an ITEV) iff the system

$$J_m(k\sqrt{n_0})A - J_m(k) = 0$$
  
$$\sqrt{n_0}J'_m(k\sqrt{n_0})A - J'_m(k) = 0$$

has a solution.

ふして 山田 ふぼやえばや 山下

that is iff 
$$d_m(k):=J_m(k\sqrt{n_0})J_m'(k)-\sqrt{n_0}J_m'(k\sqrt{n_0})J_m(k)=0$$
 .

◆□▶ ◆□▶ ◆三▶ ◆三▶ ◆□▶

that is iff 
$$d_m(k) := J_m(k\sqrt{n_0})J_m'(k) - \sqrt{n_0}J_m'(k\sqrt{n_0})J_m(k) = 0$$
 .

▲□▶ ▲□▶ ▲ 三▶ ▲ 三▶ 三三 - のへぐ

A simple calculation (based on known asymptotics of Bessel functions) yields

that is iff 
$$d_m(k) := J_m(k\sqrt{n_0})J_m'(k) - \sqrt{n_0}J_m'(k\sqrt{n_0})J_m(k) = 0$$
 .

A simple calculation (based on known asymptotics of Bessel functions) yields for  $0 < n_0 < 1$ :

$$d_m(\frac{2j\pi-\frac{\pi}{2}}{1-\sqrt{n_0}})>0 \quad \text{and} \quad d_m(\frac{2j\pi+\frac{\pi}{2}}{1-\sqrt{n_0}})<0$$

▲□▶ ▲□▶ ▲ 三▶ ▲ 三▶ 三三 - のへぐ

that is iff 
$$d_m(k) := J_m(k\sqrt{n_0})J_m'(k) - \sqrt{n_0}J_m'(k\sqrt{n_0})J_m(k) = 0$$
 .

A simple calculation (based on known asymptotics of Bessel functions) yields for  $0 < n_0 < 1$ :

$$d_m(\frac{2j\pi-\frac{\pi}{2}}{1-\sqrt{n_0}})>0 \quad \text{and} \quad d_m(\frac{2j\pi+\frac{\pi}{2}}{1-\sqrt{n_0}})<0$$

so there exists infinitely many zeroes  $\{k_j\}_{j=1}^{\infty}$  for  $d_m(\cdot)$  with

$$\frac{2j\pi - \frac{\pi}{2}}{1 - \sqrt{n_0}} < k_j < \frac{2j\pi + \frac{\pi}{2}}{1 - \sqrt{n_0}}$$

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへぐ

that is iff 
$$d_m(k) := J_m(k\sqrt{n_0})J_m'(k) - \sqrt{n_0}J_m'(k\sqrt{n_0})J_m(k) = 0$$
 .

A simple calculation (based on known asymptotics of Bessel functions) yields for  $0 < n_0 < 1$ :

$$d_m(\frac{2j\pi-\frac{\pi}{2}}{1-\sqrt{n_0}})>0 \quad \text{and} \quad d_m(\frac{2j\pi+\frac{\pi}{2}}{1-\sqrt{n_0}})<0$$

so there exists infinitely many zeroes  $\{k_j\}_{j=1}^{\infty}$  for  $d_m(\cdot)$  with

$$\frac{2j\pi - \frac{\pi}{2}}{1 - \sqrt{n_0}} < k_j < \frac{2j\pi + \frac{\pi}{2}}{1 - \sqrt{n_0}}$$

for  $1 < n_0$  :

$$d_m(\frac{2j\pi+\frac{\pi}{2}}{\sqrt{n_0}-1})>0 \quad \text{and} \quad d_m(\frac{2j\pi+\frac{3\pi}{2}}{\sqrt{n_0}-1})<0$$

so there exists infinitely many zeroes  $\{k_j\}_{j=1}^\infty$  for  $d_m(\cdot)$  with

$$\frac{2j\pi + \frac{\pi}{2}}{\sqrt{n_0 - 1}} < k_j < \frac{2j\pi + \frac{3\pi}{2}}{\sqrt{n_0 - 1}}$$

this was an example of A = B: both with countably many elements, and infinitely many non-scattering wave numbers for each density e<sup>-imθ</sup>.

▲□▶ ▲□▶ ▲ 三▶ ▲ 三▶ 三三 - のへぐ

▲□▶ ▲□▶ ▲ 三▶ ▲ 三▶ 三三 - のへぐ

We note that

$$\frac{1}{2\pi} \int_{S^1} e^{-im\theta_{\xi}} e^{ikx\cdot\xi} d\xi = \frac{1}{2\pi} \int_{-\pi}^{\pi} e^{-im\theta} e^{ik|x|\cos(\theta-\theta_x)} d\theta$$
$$= e^{-im\theta_x} \frac{1}{2\pi} \int_{-\pi}^{\pi} e^{-im\theta} e^{ik|x|\cos(\theta)} d\theta$$
$$= e^{-im\theta_x} i^m J_m(k|x|)$$

▲□▶ ▲□▶ ▲ 三▶ ▲ 三▶ 三三 - のへぐ

We note that

$$\frac{1}{2\pi} \int_{S^1} e^{-im\theta_{\xi}} e^{ikx\cdot\xi} d\xi = \frac{1}{2\pi} \int_{-\pi}^{\pi} e^{-im\theta} e^{ik|x|\cos(\theta-\theta_x)} d\theta$$
$$= e^{-im\theta_x} \frac{1}{2\pi} \int_{-\pi}^{\pi} e^{-im\theta} e^{ik|x|\cos(\theta)} d\theta$$
$$= e^{-im\theta_x} i^m J_m(k|x|)$$

▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□ ● ● ●

So  $e^{-im\theta}J_m(kr)$  is a so-called Herglotz function,

We note that

$$\frac{1}{2\pi} \int_{S^1} e^{-im\theta_{\xi}} e^{ikx\cdot\xi} d\xi = \frac{1}{2\pi} \int_{-\pi}^{\pi} e^{-im\theta} e^{ik|x|\cos(\theta-\theta_x)} d\theta$$
$$= e^{-im\theta_x} \frac{1}{2\pi} \int_{-\pi}^{\pi} e^{-im\theta} e^{ik|x|\cos(\theta)} d\theta$$
$$= e^{-im\theta_x} i^m J_m(k|x|)$$

So  $e^{-im\theta}J_m(kr)$  is a so-called Herglotz function, a function of the form  $\frac{1}{2\pi}\int_{S^1}\phi(\theta_\xi)e^{ikx\cdot\xi} d\xi$ ,

We note that

$$\frac{1}{2\pi} \int_{S^1} e^{-im\theta_{\xi}} e^{ikx\cdot\xi} d\xi = \frac{1}{2\pi} \int_{-\pi}^{\pi} e^{-im\theta} e^{ik|x|\cos(\theta-\theta_x)} d\theta$$
$$= e^{-im\theta_x} \frac{1}{2\pi} \int_{-\pi}^{\pi} e^{-im\theta} e^{ik|x|\cos(\theta)} d\theta$$
$$= e^{-im\theta_x} i^m J_m(k|x|)$$

So  $e^{-im\theta}J_m(kr)$  is a so-called Herglotz function, a function of the form  $\frac{1}{2\pi}\int_{S^1}\phi(\theta_\xi)e^{ikx\cdot\xi} d\xi$ , with density  $\phi(\theta) = (-i)^m e^{-im\theta}$ .

・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・

We note that

$$\frac{1}{2\pi} \int_{S^1} e^{-im\theta_{\xi}} e^{ikx\cdot\xi} d\xi = \frac{1}{2\pi} \int_{-\pi}^{\pi} e^{-im\theta} e^{ik|x|\cos(\theta-\theta_x)} d\theta$$
$$= e^{-im\theta_x} \frac{1}{2\pi} \int_{-\pi}^{\pi} e^{-im\theta} e^{ik|x|\cos(\theta)} d\theta$$
$$= e^{-im\theta_x} i^m J_m(k|x|)$$

So  $e^{-im\theta}J_m(kr)$  is a so-called Herglotz function, a function of the form  $\frac{1}{2\pi}\int_{S^1}\phi(\theta_\xi)e^{ikx\cdot\xi} d\xi$ , with density  $\phi(\theta) = (-i)^m e^{-im\theta}$ .

・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・

So what happens if you perturb the disk to an ellipse?

Introduction

#### Finiteness of Non-scattering Wave Numbers

Connection with Schiffer's conjecture



Theorem (Herglotz Waves [Vogelius-Xiao '20]) Let  $\phi \in C^{\infty}(\mathbb{S}^1) \setminus \{0\}$  and  $n_0 \in \mathbb{R}_+ \setminus \{1\}$ . Let  $\Omega$  denote the ellipse

$$\Omega = \Omega_{A,B} = \{ (x_1, x_2) : \frac{x_1^2}{A^2} + \frac{x_2^2}{B^2} < 1 \}$$

▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□ ● ● ●

Theorem (Herglotz Waves [Vogelius-Xiao '20]) Let  $\phi \in C^{\infty}(\mathbb{S}^1) \setminus \{0\}$  and  $n_0 \in \mathbb{R}_+ \setminus \{1\}$ . Let  $\Omega$  denote the ellipse

$$\Omega = \Omega_{A,B} = \{ (x_1, x_2) : \frac{x_1^2}{A^2} + \frac{x_2^2}{B^2} < 1 \}$$

$$u_{k,in} = \mathbf{H}[k,\phi](x) = \frac{1}{2\pi} \int_{-\pi}^{\pi} \phi(\theta) e^{ik\xi(\theta) \cdot x} d\theta .$$

◆□▶ ◆□▶ ◆ 臣▶ ◆ 臣▶ 三臣 - ∽ � � �

Theorem (Herglotz Waves [Vogelius-Xiao '20]) Let  $\phi \in C^{\infty}(\mathbb{S}^1) \setminus \{0\}$  and  $n_0 \in \mathbb{R}_+ \setminus \{1\}$ . Let  $\Omega$  denote the ellipse

$$\Omega = \Omega_{A,B} = \{ (x_1, x_2) : \frac{x_1^2}{A^2} + \frac{x_2^2}{B^2} < 1 \}$$

There exists a positive constant  $\epsilon$  (independent of  $\phi$ ), so that if  $0 < |\frac{B}{A} - 1| < \epsilon$  then one can find at most finitely many positive wave numbers  $k = k_j, j = 1, 2, ..., N$ , (possibly depending on  $\phi$ ) for which the problem

$$\begin{cases} \Delta u_{k,sc} + k^2 n_0 u_{k,sc} = k^2 (1 - n_0) u_{k,in} & \text{in } \Omega, \\ u_{k,sc} = \partial_{\nu} u_{k,sc} = 0 & \text{on } \partial\Omega \end{cases}$$

admits a solution  $u_{k,sc}$ , with

$$u_{k,in} = \mathbf{H}[k,\phi](x) = \frac{1}{2\pi} \int_{-\pi}^{\pi} \phi(\theta) e^{ik\xi(\theta) \cdot x} d\theta .$$

・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・

### $\mathsf{Cont'd}$

Consequence: At most finitely many non-scattering wave numbers for (Ω, n<sub>0</sub>, H[k, φ]).

Consequence: At most finitely many non-scattering wave numbers for (Ω, n<sub>0</sub>, H[k, φ]).

▲□▶ ▲□▶ ▲ 三▶ ▲ 三▶ 三 のへぐ

Some generalizations ([Vogelius-Xiao, in preparation]):

Consequence: At most finitely many non-scattering wave numbers for (Ω, n<sub>0</sub>, H[k, φ]).

Some generalizations ([Vogelius-Xiao, in preparation]):

Ellipses with

$$\frac{1}{1+\sqrt{n_0}} < B^2/A^2 < 1+\sqrt{n_0}, \qquad A \neq B.$$

▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□ ● ● ●

Consequence: At most finitely many non-scattering wave numbers for (Ω, n<sub>0</sub>, H[k, φ]).

Some generalizations ([Vogelius-Xiao, in preparation]):

Ellipses with

$$\frac{1}{1+\sqrt{n_0}} < B^2/A^2 < 1+\sqrt{n_0}, \qquad A \neq B.$$

▲□▶ ▲□▶ ▲□▶ ▲□▶ ■ ●の00

Small  $C^2$  perturbations of ellipses.

- Consequence: At most finitely many non-scattering wave numbers for (Ω, n<sub>0</sub>, H[k, φ]).
- Some generalizations ([Vogelius-Xiao, in preparation]):

Ellipses with

$$\frac{1}{1+\sqrt{n_0}} < B^2/A^2 < 1+\sqrt{n_0}, \qquad A \neq B.$$

▲□▶ ▲□▶ ▲□▶ ▲□▶ ■ ●の00

Small C<sup>2</sup> perturbations of ellipses.
 Small C<sup>2</sup> perturbations of disks.

- Consequence: At most finitely many non-scattering wave numbers for (Ω, n<sub>0</sub>, H[k, φ]).
- Some generalizations ([Vogelius-Xiao, in preparation]):

Ellipses with

$$\frac{1}{1+\sqrt{n_0}} < B^2/A^2 < 1+\sqrt{n_0}, \qquad A \neq B.$$

▲□▶ ▲□▶ ▲□▶ ▲□▶ ■ ●の00

Small C<sup>2</sup> perturbations of ellipses.
 Small C<sup>2</sup> perturbations of disks.

- Consequence: At most finitely many non-scattering wave numbers for (Ω, n<sub>0</sub>, H[k, φ]).
- Some generalizations ([Vogelius-Xiao, in preparation]):

Ellipses with

$$\frac{1}{1+\sqrt{n_0}} < B^2/A^2 < 1+\sqrt{n_0}, \qquad A \neq B.$$

Small C<sup>2</sup> perturbations of ellipses.
 Small C<sup>2</sup> perturbations of disks.

What happens for other incident waves, e.g., plane waves  $e^{ikx\cdot\xi}$ ?

Suppose Ω is C<sup>1,α</sup>, and suppose there exists a non-scattering wave number for (Ω, n<sub>0</sub>, e<sup>ikx·ξ</sup>). Then Ω must be real analytic (in other words its boundary is a real analytic hypersurface).

- If Ω is a ball then there are exists no non-scattering wave number for (Ω, n<sub>0</sub>, e<sup>ikx·ξ</sup>).
- This raises a natural question:

- Suppose Ω is C<sup>1,α</sup>, and suppose there exists a non-scattering wave number for (Ω, n<sub>0</sub>, e<sup>ikx.ξ</sup>). Then Ω must be real analytic (in other words its boundary is a real analytic hypersurface).
- If Ω is a ball then there are exists no non-scattering wave number for (Ω, n<sub>0</sub>, e<sup>ikx·ξ</sup>).

This raises a natural question: if Ω is real analytic but not a ball what might one say about the number of non-scattering wave numbers for (Ω, n<sub>0</sub>, e<sup>ikx·ξ</sup>)

#### Theorem (Plane Waves [Vogelius-Xiao '20])

Let  $\Omega$  be a bounded, strictly convex  $C^{2,\alpha}$  domain in  $\mathbb{R}^2$ , and suppose  $n_0 \in \mathbb{R}_+ \setminus \{1\}$ . Given  $u_{k,in} = e^{ikx \cdot \xi}$ , for a fixed direction  $\xi \in \mathbb{S}^1$ , there exist at most finitely many positive wave numbers  $k = k_j$ ,  $j = 1, \ldots, N$ , such that

$$\begin{cases} \Delta u_{k,sc} + k^2 n_0 u_{k,sc} = k^2 (1 - n_0) u_{k,in} & \text{in } \Omega, \\ u_{k,sc} = \partial_{\nu} u_{k,sc} = 0 & \text{on } \partial \Omega, \end{cases}$$

A D N A 目 N A E N A E N A B N A C N

admits a solution  $u_{k,sc}$ .

#### Theorem (Plane Waves [Vogelius-Xiao '20])

Let  $\Omega$  be a bounded, strictly convex  $C^{2,\alpha}$  domain in  $\mathbb{R}^2$ , and suppose  $n_0 \in \mathbb{R}_+ \setminus \{1\}$ . Given  $u_{k,in} = e^{ikx \cdot \xi}$ , for a fixed direction  $\xi \in \mathbb{S}^1$ , there exist at most finitely many positive wave numbers  $k = k_j$ ,  $j = 1, \ldots, N$ , such that

$$\begin{cases} \Delta u_{k,sc} + k^2 n_0 u_{k,sc} = k^2 (1 - n_0) u_{k,in} & \text{in } \Omega, \\ u_{k,sc} = \partial_{\nu} u_{k,sc} = 0 & \text{on } \partial \Omega, \end{cases}$$

admits a solution  $u_{k,sc}$ .

 Consequence: At most finitely many non-scattering wave numbers for (Ω, n<sub>0</sub>, e<sup>ikx·ξ</sup>).

▲□▶ ▲□▶ ▲□▶ ▲□▶ ■ ● ●

Suppose there exists a solution  $u_{k,sc}$  to

Suppose there exists a solution  $u_{k,sc}$  to

$$\begin{split} \Delta u_{k,in} + k^2 u_{k,in} &= 0 & \text{in } \Omega, \\ \Delta u_{k,sc} + k^2 n_0 u_{k,sc} &= k^2 (1 - n_0) u_{k,in} & \\ u_{k,sc} &= \partial_{\nu} u_{k,sc} &= 0 & \text{on } \partial \Omega. \end{split}$$

Suppose there exists a solution  $u_{k,sc}$  to

$$\begin{split} \Delta u_{k,in} + k^2 u_{k,in} &= 0 & \text{in } \Omega, \\ \Delta u_{k,sc} + k^2 n_0 u_{k,sc} &= k^2 (1 - n_0) u_{k,in} & \\ u_{k,sc} &= \partial_\nu u_{k,sc} &= 0 & \text{on } \partial \Omega. \end{split}$$

then for any solution w to  $\Delta w + k^2 n_0 w = 0$ ,

$$k^2 \int_{\Omega} (1 - n_0) \, u_{k,in} \, w \, dx = \int_{\partial \Omega} w \, \partial_{\nu} u_{k,sc} - u_{k,sc} \, \partial_{\nu} w \, d\sigma(x) = 0.$$

Suppose there exists a solution  $u_{k,sc}$  to

$$\begin{split} \Delta u_{k,in} + k^2 u_{k,in} &= 0 & \text{in } \Omega, \\ \Delta u_{k,sc} + k^2 n_0 u_{k,sc} &= k^2 (1 - n_0) u_{k,in} & \\ u_{k,sc} &= \partial_\nu u_{k,sc} &= 0 & \text{on } \partial \Omega. \end{split}$$

then for any solution w to  $\Delta w + k^2 n_0 w = 0$ ,

$$k^2 \int_{\Omega} (1 - n_0) \, u_{k,in} \, w \, dx = \int_{\partial \Omega} w \, \partial_{\nu} u_{k,sc} - u_{k,sc} \, \partial_{\nu} w \, d\sigma(x) = 0.$$

We may in particular take  $w=e^{ik\sqrt{n_0}\,\eta\cdot x}$  for any  $\eta\in\mathbb{S}^1$ 

$$\int_{\Omega} (1 - n_0) \, u_{k,in} \, e^{ik\sqrt{n_0} \, \eta \cdot x} \, dx = 0$$

(ロ)、(型)、(E)、(E)、 E) の(()

Approach:

$$\int_{\Omega} (1 - n_0) \, u_{k,in} \, e^{ik\sqrt{n_0} \, \eta \cdot x} \, dx = 0$$

 if the are are infinitely many non-scattering wave numbers then the above identity holds for infinitely many k (for all η).

▲□▶▲□▶▲≡▶▲≡▶ ≡ めぬぐ

$$\int_{\Omega} (1 - n_0) \, u_{k,in} \, e^{ik\sqrt{n_0} \, \eta \cdot x} \, dx = 0$$

 if the are are infinitely many non-scattering wave numbers then the above identity holds for infinitely many k (for all η).

▲□▶▲□▶▲≡▶▲≡▶ ≡ めぬぐ

$$\int_{\Omega} (1 - n_0) \, u_{k,in} \, e^{ik\sqrt{n_0} \, \eta \cdot x} \, dx = 0$$

 if the are infinitely many non-scattering wave numbers then the above identity holds for infinitely many k (for all η). The k's must be countable and tend to ∞ (they are transmission eigenvalues).

▲□▶ ▲□▶ ▲□▶ ▲□▶ ■ ●の00

$$\int_{\Omega} (1 - n_0) \, u_{k,in} \, e^{ik\sqrt{n_0} \, \eta \cdot x} \, dx = 0$$

- if the are infinitely many non-scattering wave numbers then the above identity holds for infinitely many k (for all η). The k's must be countable and tend to ∞ (they are transmission eigenvalues).
- The left hand side is an oscillatory integral, the asymptotics of which we may derive by the method of stationary phase.

▲□▶ ▲□▶ ▲□▶ ▲□▶ ■ ●の00

$$\int_{\Omega} (1 - n_0) \, u_{k,in} \, e^{ik\sqrt{n_0} \, \eta \cdot x} \, dx = 0$$

- if the are infinitely many non-scattering wave numbers then the above identity holds for infinitely many k (for all η). The k's must be countable and tend to ∞ (they are transmission eigenvalues).
- The left hand side is an oscillatory integral, the asymptotics of which we may derive by the method of stationary phase.

▲□▶ ▲□▶ ▲□▶ ▲□▶ ■ ●の00

$$\int_{\Omega} (1 - n_0) \, u_{k,in} \, e^{ik\sqrt{n_0} \, \eta \cdot x} \, dx = 0$$

- if the are infinitely many non-scattering wave numbers then the above identity holds for infinitely many k (for all η). The k's must be countable and tend to ∞ (they are transmission eigenvalues).
- The left hand side is an oscillatory integral, the asymptotics of which we may derive by the method of stationary phase.

As a result we obtain a contradiction when

$$\int_{\Omega} (1 - n_0) \, u_{k,in} \, e^{ik\sqrt{n_0} \, \eta \cdot x} \, dx = 0$$

- if the are infinitely many non-scattering wave numbers then the above identity holds for infinitely many k (for all η). The k's must be countable and tend to ∞ (they are transmission eigenvalues).
- The left hand side is an oscillatory integral, the asymptotics of which we may derive by the method of stationary phase.

As a result we obtain a contradiction when

$$\int_{\Omega} (1 - n_0) \, u_{k,in} \, e^{ik\sqrt{n_0} \, \eta \cdot x} \, dx = 0$$

- if the are infinitely many non-scattering wave numbers then the above identity holds for infinitely many k (for all η). The k's must be countable and tend to ∞ (they are transmission eigenvalues).
- The left hand side is an oscillatory integral, the asymptotics of which we may derive by the method of stationary phase.
- As a result we obtain a contradiction when (1) u<sub>k,in</sub> is a plane wave or

$$\int_{\Omega} (1 - n_0) \, u_{k,in} \, e^{ik\sqrt{n_0} \, \eta \cdot x} \, dx = 0$$

- if the are are infinitely many non-scattering wave numbers then the above identity holds for infinitely many k (for all η). The k's must be countable and tend to ∞ (they are transmission eigenvalues).
- The left hand side is an oscillatory integral, the asymptotics of which we may derive by the method of stationary phase.
- As a result we obtain a contradiction when (1) u<sub>k,in</sub> is a plane wave or (2) when u<sub>k,in</sub> is a Herglotz wave and Ω is a proper ellipse.

Introduction

Finiteness of Non-scattering Wave Numbers

Connection with Schiffer's conjecture

◆□ ▶ ◆□ ▶ ◆ 臣 ▶ ◆ 臣 ▶ ○ 臣 ○ のへで

$$\begin{split} \Omega \text{ always scatters if:} \\ \Delta u_{k,in} + k^2 u_{k,in} &= 0 \\ \begin{cases} \Delta u_{k,sc} + k^2 n_0 u_{k,sc} &= k^2 (1-n_0) u_{k,in} & \text{ in } \Omega, \\ u_{k,sc} &= \partial_\nu u_{k,sc} &= 0 & \text{ on } \partial \Omega. \end{cases} \end{split}$$

(ロ)、(型)、(E)、(E)、 E) の(()

has **no** solution  $u_{k,sc}$  for any wave number k.

$$\begin{split} \Omega \text{ always scatters if:} \\ \Delta u_{k,in} + k^2 u_{k,in} &= 0 \\ \begin{cases} \Delta u_{k,sc} + k^2 n_0 u_{k,sc} = k^2 (1-n_0) u_{k,in} & \text{ in } \Omega, \\ u_{k,sc} &= \partial_\nu u_{k,sc} = 0 & \text{ on } \partial \Omega. \end{cases} \end{split}$$

has **no** solution  $u_{k,sc}$  for any wave number k.

We shall say a domain has the Schiffer Property if the problem

$$\begin{split} \Omega \text{ always scatters if:} \\ \Delta u_{k,in} + k^2 u_{k,in} &= 0 \\ \begin{cases} \Delta u_{k,sc} + k^2 n_0 u_{k,sc} &= k^2 (1-n_0) u_{k,in} & \text{ in } \Omega, \\ u_{k,sc} &= \partial_\nu u_{k,sc} &= 0 & \text{ on } \partial \Omega. \end{cases} \end{split}$$

has **no** solution  $u_{k,sc}$  for any wave number k.

We shall say a domain has the Schiffer Property if the problem

$$\begin{cases} \Delta v + \lambda v = 1 & \text{in } \Omega, \\ v = \partial_{\nu} v = 0 & \text{on } \partial \Omega. \end{cases}$$

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三回 のへぐ

$$\begin{split} \Omega \text{ always scatters if:} \\ \Delta u_{k,in} + k^2 u_{k,in} &= 0 \\ \begin{cases} \Delta u_{k,sc} + k^2 n_0 u_{k,sc} &= k^2 (1-n_0) u_{k,in} & \text{ in } \Omega, \\ u_{k,sc} &= \partial_\nu u_{k,sc} &= 0 & \text{ on } \partial \Omega. \end{cases} \end{split}$$

has **no** solution  $u_{k,sc}$  for any wave number k.

We shall say a domain has the Schiffer Property if the problem

$$\begin{cases} \Delta v + \lambda v = 1 & \text{ in } \Omega, \\ v = \partial_{\nu} v = 0 & \text{ on } \partial \Omega. \end{cases}$$

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへぐ

has **no** solution for any  $\lambda$ .

$$\begin{split} \Omega \text{ always scatters if:} \\ \Delta u_{k,in} + k^2 u_{k,in} &= 0 \\ \begin{cases} \Delta u_{k,sc} + k^2 n_0 u_{k,sc} = k^2 (1-n_0) u_{k,in} & \text{ in } \Omega, \\ u_{k,sc} &= \partial_\nu u_{k,sc} = 0 & \text{ on } \partial \Omega. \end{cases} \end{split}$$

has **no** solution  $u_{k,sc}$  for any wave number k.

We shall say a domain has the Schiffer Property if the problem

$$\begin{cases} \Delta v + \lambda v = 1 & \text{ in } \Omega, \\ v = \partial_{\nu} v = 0 & \text{ on } \partial \Omega. \end{cases}$$

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへぐ

has **no** solution for any  $\lambda$ .

This is equivalent to

$$\begin{cases} \Delta w + \lambda w = 0 & \text{ in } \Omega, \\ \partial_{\nu} w = 0, \quad w = const, & \text{ on } \partial \Omega. \end{cases}$$

has **no** non-trivial solution for  $\lambda \neq 0$ 

$$\begin{cases} \Delta w + \lambda w = 0 & \text{ in } \Omega, \\ \partial_{\nu} w = 0, \quad w = \textit{const}, & \text{ on } \partial \Omega. \end{cases}$$

has **no** non-trivial solution for  $\lambda \neq 0$  or alternatively: there are no non-trivial Neumann eigenfunctions (for the Laplacian,  $\Delta$ ) with constant Dirichlet boundary trace.

▲□▶ ▲□▶ ▲□▶ ▲□▶ ■ ●の00

$$\begin{cases} \Delta w + \lambda w = 0 & \text{ in } \Omega, \\ \partial_{\nu} w = 0, \quad w = const, & \text{ on } \partial \Omega. \end{cases}$$

has **no** non-trivial solution for  $\lambda \neq 0$  or alternatively: there are no non-trivial Neumann eigenfunctions (for the Laplacian,  $\Delta$ ) with constant Dirichlet boundary trace.

Conjecture: The only simply connected domains in  $\mathbb{R}^d$  that fail to have the Schiffer Property are balls.

・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・

$$\begin{cases} \Delta w + \lambda w = 0 & \text{ in } \Omega, \\ \partial_{\nu} w = 0, \quad w = const, & \text{ on } \partial \Omega. \end{cases}$$

has **no** non-trivial solution for  $\lambda \neq 0$  or alternatively: there are no non-trivial Neumann eigenfunctions (for the Laplacian,  $\Delta$ ) with constant Dirichlet boundary trace.

Conjecture: The only simply connected domains in  $\mathbb{R}^d$  that fail to have the Schiffer Property are balls.

Note that: balls have infinitely many Neumann eigenvalues with corresponding eigenfunctions of constant Dirichlet boundary trace.

・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・

Some interesting (and related) results:

<ロト < 団ト < 団ト < 団ト < 団ト 三 のQの</p>

Some interesting (and related) results:

 Any Lipschitz domain Ω that fails to have the Schiffer Property, is real analytic. [Williams '81]

▲□▶ ▲□▶ ▲ 三▶ ▲ 三▶ 三三 - のへぐ

Some interesting (and related) results:

- Any Lipschitz domain Ω that fails to have the Schiffer Property, is real analytic. [Williams '81]
- If Ω is a simply connected Lipschitz domain with infinitely many Neumann eigenvalues with corresponding eigenfunctions of constant Dirichlet trace, then Ω is a ball. [Berenstein-Yang '87, Vogelius '94]